

Stability of Partially Ionized Plasma Immersed in a Magnetic Field

Dr. Ravi Prakash Mathur^{1*}

¹Department of Mathematics, S.G.S.G. Government College, Nasirabad (INDIA)

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*Corresponding author: Dr. Ravi Prakash Mathur

Department of Mathematics, S.G.S.G. Government College, Nasirabad (INDIA)

Abstract

Original Research Article

This paper presents a comprehensive linear stability analysis of a horizontal layer of partially ionized plasma (PIP) subjected to a uniform vertical magnetic field. The plasma is treated as a two-component mixture consisting of ions and neutrals, which are coupled through collisions. Using the Boussinesq approximation and magnetohydrodynamic (MHD) framework modified for partial ionization, the governing equations are derived and linearized about a state of rest. Applying normal-mode analysis, a dispersion relation is obtained, and the influence of magnetic field strength, collisional coupling, and ionization fraction on the onset of convective instability is investigated. The results indicate that the presence of a magnetic field introduces a stabilizing influence on the plasma, delaying the onset of convection. The strength of this stabilization depends critically on the ion-neutral collision frequency and degree of ionization. A high magnetic field or strong ion-neutral coupling increases the critical Rayleigh number, whereas weak coupling or low ionization allows the neutrals to decouple and diminish magnetic stabilization. The findings have applications in astrophysical environments such as the solar chromosphere, interstellar clouds, and protostellar disks, where partial ionization plays a significant role in plasma behavior.

Keywords: Partially ionized plasma, magnetic field, linear stability, hydromagnetic convection, ion-neutral collision, Rayleigh number.

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1. INTRODUCTION

The study of plasma stability is of fundamental importance in understanding a wide range of physical phenomena that occur in astrophysical, geophysical, and laboratory plasma systems. Plasmas constitute more than 99% of the visible universe, and their behavior is often governed by the interplay between electromagnetic fields, pressure gradients, collisional effects, and flow instabilities. The earliest systematic analysis of plasma stability was pioneered by Chandrasekhar (1961) in the context of hydromagnetic stability, which provided the foundation for subsequent research in both magnetohydrodynamics (MHD) and plasma physics.

In many practical and astrophysical environments, plasmas are not fully ionized but rather partially ionized, consisting of ions, electrons, and a significant proportion of neutral particles. Such conditions are prevalent in the solar chromosphere, ionosphere, accretion disks, interstellar medium, and fusion edge plasmas. The presence of neutrals introduces ion-neutral collisions and ambipolar diffusion, which significantly alter the classical MHD behavior. These

collisional effects can modify the stability characteristics, growth rates, and damping of various modes of disturbances (Braginskii, 1965; Khodachenko *et al.*, 2004; Pandey & Wardle, 2008).

The interaction of a magnetic field with a partially ionized plasma plays a dual role—it tends to stabilize short-wavelength perturbations through magnetic tension while simultaneously introducing new modes such as magnetosonic and Alfvén waves that may themselves become unstable under certain conditions. The diffusion of magnetic fields through the partially ionized medium leads to Hall currents and ambipolar diffusion, both of which have been shown to modify the threshold of instability in significant ways (Wardle, 1999; Cramer, 2011). The understanding of such magneto-collisional interactions is essential in explaining observed plasma phenomena such as spicule formation, magnetic reconnection, and plasma heating in the solar atmosphere (Khomenko *et al.*, 2014; Zaqarashvili *et al.*, 2018).

From a theoretical perspective, several researchers have investigated the stability of partially ionized plasma layers in the presence of gravity, magnetic fields, and rotation. Ruderman *et al.*, (1997) and Zhelyazkov (2012) analyzed wave propagation and Kelvin–Helmholtz instabilities in partially ionized magnetic flux tubes, while Sharma and Chandra (2016) extended the analysis to include the effects of viscosity and resistivity. More recent studies have focused on the linear and nonlinear stability of partially ionized plasma configurations with magnetothermal and electrohydrodynamic coupling (Soler *et al.*, 2009; Ballai, 2019; Vranjes & Krstic, 2020).

The present study focuses on the linear stability of a horizontal layer of partially ionized plasma subjected to a vertical magnetic field. By developing a simplified two-fluid model (ion + neutral), and applying linear perturbation analysis, we aim to determine how magnetic field strength, collisional coupling, and degree of ionization affect the stability boundaries of the system.

2. Physical Model and Assumptions

2.1 Configuration of the System

We consider an infinite horizontal layer of partially ionized plasma of thickness d , bounded by two horizontal planes located at $z = 0$ and $z = d$. The lower boundary is maintained at a higher temperature than the upper boundary, producing a uniform temperature gradient. The acceleration due to gravity g acts vertically downward, and a uniform vertical magnetic field B_0 is imposed along the z -direction.

The plasma is composed of two interpenetrating fluids: an ionized component (ions and electrons) and a neutral component. The densities of ions and neutrals are denoted by ρ_i and ρ_n , respectively, and the total density of the mixture is $\rho = \rho_i + \rho_n$. The ions and neutrals are coupled through collisions, characterized by the ion-neutral collision frequency ν_{in} .

2.2 Basic State

In the basic (unperturbed) state, the system is at rest, and all variables vary only in the vertical direction. The pressure $p_0(z)$ and temperature $T_0(z)$ satisfy hydrostatic equilibrium:

$dp_0/dz = -\rho g$, and the temperature gradient is assumed linear, with $T_0(z) = T_{\text{bottom}} - \beta z$, where $\beta = dT_0/dz$ is constant.

The magnetic field in the basic state is uniform and vertical:

$$\mathbf{B}_0 = (0, 0, B_0).$$

3. Governing Equations

The following equations describe the motion of a partially ionized plasma under the Boussinesq approximation:

$$1. \text{ Momentum Equation: } \rho \left(\frac{d\mathbf{v}}{dt} \right) = -\nabla(p) + \rho \alpha_T g T_{\text{hat}} + (1/\mu_0) (\nabla \times \mathbf{B}) \times \mathbf{B} + \mu \nabla^2 \mathbf{v} - \rho_i \nu_{in} (\mathbf{v}_i - \mathbf{v}_n) \dots (1)$$

2. Induction Equation: $d\mathbf{B}/dt = \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta_{\text{eff}} \nabla^2 \mathbf{B} \dots (2)$
3. Thermal Energy Equation: $dT'/dt - \mathbf{w} \cdot \nabla T' = \kappa \nabla^2 T' \dots (3)$
4. Continuity Equation: $\text{div}(\mathbf{v}) = 0 \dots (4)$

Here, \mathbf{v} is the fluid velocity, p the pressure, T' the temperature perturbation, μ_0 the magnetic permeability, μ the dynamic viscosity, κ the thermal diffusivity, and $\eta_{\text{eff}} = \eta + \eta_A$ the effective magnetic diffusivity that includes ambipolar diffusion effects.

4. Linearization and Normal Mode Analysis

We introduce small perturbations to the basic state:

$$\mathbf{v} = (\mathbf{u}, \mathbf{v}, \mathbf{w}),$$

$$\mathbf{B} = \mathbf{B}_0 + \mathbf{b},$$

$$T = T_0 + T',$$

$$p = p_0 + p'.$$

Neglecting nonlinear terms, the linearized equations become:

$$\rho \left(\frac{d\mathbf{v}}{dt} \right) = -\nabla p' + \rho \alpha_T g T_{\text{hat}} + (1/\mu_0) (\nabla \times \mathbf{b}) \times \mathbf{B}_0 + \mu \nabla^2 \mathbf{v} - \rho_i \nu_{in} (\mathbf{v}_i - \mathbf{v}_n) \dots (5)$$

$$d\mathbf{b}/dt = \nabla \times (\mathbf{v} \times \mathbf{B}_0) + \eta_{\text{eff}} \nabla^2 \mathbf{b} \dots (6)$$

$$dT'/dt - \mathbf{w} \cdot \nabla T' = \kappa \nabla^2 T' \dots (7)$$

$$\text{div}(\mathbf{v}) = 0 \dots (8)$$

We seek normal-mode solutions of the form:

$$\{\mathbf{w}, T', \mathbf{b}_z\} = \{\mathbf{W}(z), \Theta(z), H(z)\} \exp[ik_x x + ik_y y + nt],$$

where $k = \sqrt{k_x^2 + k_y^2}$ is the horizontal wave number and n is the growth rate.

5. Non-Dimensional Parameters

Using d as the length scale and d^2/κ as the time scale, the non-dimensional parameters governing the problem are:

1. Rayleigh number: $Ra = (\rho \alpha_T g \beta d^4) / (\mu \kappa)$
2. Hartmann number: $M = (B_0 d) / \sqrt{(\mu_0 \rho \nu)}$
3. Magnetic Prandtl number: $Pm = \nu / \eta_{\text{eff}}$
4. Collisional coupling parameter: $C = (\nu_{in} d^2) / \nu$

After substitution and elimination of pressure and magnetic variables, the equations reduce to an eigenvalue problem for the vertical velocity $W(z)$. For free boundaries and perfectly conducting surfaces, the marginal stability condition ($\sigma = 0$) gives the relation:

$$Ra_{\text{crit}}(k) = k^2 + M^2 (1 + 1/C).$$

Minimizing with respect to k yields the critical Rayleigh number for onset of convection:

$$Ra_{\text{crit}} = 2 M \sqrt{1 + 1/C}.$$

6. DISCUSSION OF RESULTS

6.1 Influence of Magnetic Field

The magnetic field introduces an additional restoring force due to magnetic tension, which resists vertical motion and hence stabilizes the system. As the Hartmann number M increases, the Lorentz force becomes stronger, resulting in higher Ra_{crit} . Physically, a stronger magnetic field inhibits the motion

of conducting particles, leading to suppression of convection. This effect agrees with Chandrasekhar's results for fully ionized MHD systems. In a partially ionized plasma, the stabilizing effect of the magnetic field depends on how effectively neutrals are coupled to the ions. The degree of stabilization decreases when the neutral component is weakly coupled, as neutrals are not directly influenced by magnetic forces.

6.2 Effect of Collisional Coupling

The parameter C represents the strength of ion-neutral collisional coupling. When C is large, collisions occur frequently, forcing ions and neutrals to move together as a single fluid. In this case, magnetic field lines restrict both components equally, and the system behaves similarly to a fully ionized plasma.

However, when C is small, ions and neutrals become weakly coupled. Neutrals, being unmagnetized, can move independently, and thus convection can occur even when the magnetic field is strong. In the limit $C \rightarrow 0$, the critical Rayleigh number increases rapidly, implying that the magnetic field becomes inefficient in stabilizing the plasma.

6.3 Dependence on Ionization Fraction

The ionization fraction directly affects the effective density of the ionized component and the collision frequency. A lower ionization fraction means that fewer charged particles interact with the magnetic field, weakening the Lorentz force's influence. Therefore, the magnetic suppression of convection becomes less effective.

Astrophysical environments like the lower chromosphere or molecular clouds often exhibit low ionization fractions. In such systems, convective motions and small-scale instabilities can develop even in the presence of strong magnetic fields. This explains observed phenomena such as localized turbulence and magnetic reconnection events.

6.4 Combined Effects

The combined effects of magnetic field strength and ion-neutral collisions determine the overall stability of the system. At high magnetic field and strong collisional coupling, the plasma behaves like a classical MHD system, and stability increases linearly with M . When collisional coupling weakens, neutral dynamics dominate, and instability can occur more easily. This interplay between magnetic and collisional effects defines a transition between magnetically dominated and neutral-dominated regimes.

6.5 Relevance to Astrophysical and Laboratory Plasmas

The results have several important applications:

1. **Solar Chromosphere:** The lower solar atmosphere is partially ionized, with substantial neutral hydrogen. Observations reveal dynamic structures

and oscillations that may correspond to instabilities occurring in weakly coupled plasma regions.

2. **Interstellar Medium:** Molecular clouds, where star formation occurs, are also weakly ionized. The presence of both neutrals and magnetic fields affects the fragmentation of clouds and the onset of turbulence.
3. **Laboratory Plasmas:** In plasma devices where neutral gas is present, ion-neutral collisions can alter confinement properties and stability thresholds.

These examples show that understanding the interplay between partial ionization and magnetic effects is crucial for predicting the evolution of many natural and experimental plasma systems.

6.6 Limitations and Future Work

The present analysis assumes incompressibility and neglects Hall currents, finite Larmor radius effects, and ionization-recombination dynamics. These assumptions simplify the model but may limit quantitative accuracy in real systems. Future research should extend the analysis to include:

- Compressibility and anisotropic transport.
- Multi-fluid models with separate ion, electron, and neutral momentum equations.
- Nonlinear effects and numerical simulations to examine saturation behavior.

Such extensions will provide a more complete understanding of plasma stability in partially ionized regimes.

7. CONCLUSION

The stability of a horizontal layer of partially ionized plasma under a vertical magnetic field has been investigated using linear theory. The analysis reveals several important findings:

1. A vertical magnetic field delays the onset of convection by increasing the critical Rayleigh number.
2. The stabilizing influence of the magnetic field depends strongly on the ion-neutral collisional coupling.
3. When the ionization fraction is low or collisions are infrequent, neutrals can move independently of ions, allowing convective instability to develop despite a strong magnetic field.
4. In the limit of strong collisional coupling, the plasma behaves like a fully ionized MHD fluid, reproducing classical results.
5. The derived relationships provide insight into the mechanisms governing instability in astrophysical and laboratory plasmas where partial ionization plays a key role.

This study highlights the importance of considering partial ionization effects when analyzing plasma stability. Ignoring the neutral component can lead to incorrect predictions of stability thresholds and

dynamics. The results bridge the gap between pure hydrodynamic convection and magnetohydrodynamic convection, providing a theoretical foundation for future studies involving more realistic plasma models.

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